

## Chapter 5: Angular Momentum

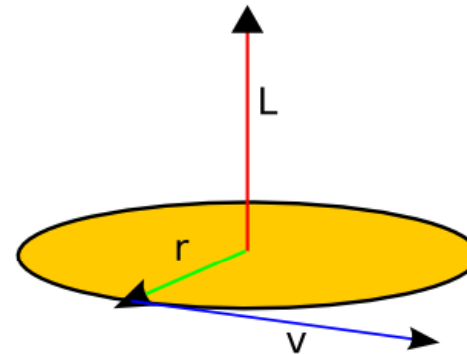
### What is angular momentum?

- A quantity used in describing rotating systems.

- It is defined as a cross product of linear momentum of the particle ( $p=mV$ ) and a position vector of a particle ( $r$ ).

### Why do we care about angular momentum?

- To describe the motion of electron around nucleus



$$\vec{L} = \vec{r} \times \vec{p}$$



*Conservation of angular momentum*

### Review of vectors

- We use cartesian coordinates to represent vectors:

$$\vec{A} = A_x \cdot \vec{i} + A_y \cdot \vec{j} + A_z \cdot \vec{k}$$

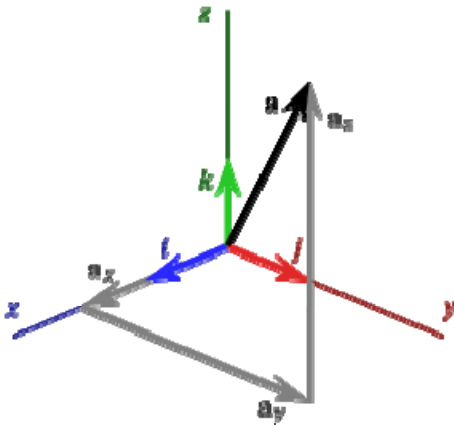
The magnitude of the vector is its length:

$$|\vec{A}| = \sqrt{A_x^2 + A_y^2 + A_z^2}$$

Sum of two vectors:  $\vec{A} + \vec{B} = (A_x + B_x) \vec{i} + (A_y + B_y) \vec{j} + (A_z + B_z) \vec{k}$

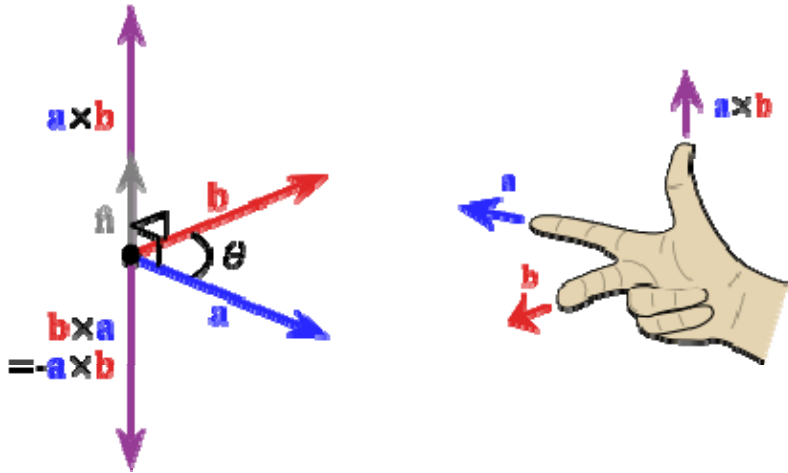
Dot (scalar) product:  $\vec{A} \cdot \vec{B} = |\vec{A}| \cdot |\vec{B}| \cdot \cos \theta = A_x \cdot B_x + A_y \cdot B_y + A_z \cdot B_z$

scalar  
angle between the vectors



-The cross (vector) product: it is a vector with magnitude:  $|\vec{A} \times \vec{B}| = |\vec{A}| \cdot |\vec{B}| \cdot \sin \theta$

Its direction is perpendicular to the plane defined by vectors A and B:



How to calculate vector product:

$$\vec{A} \times \vec{B} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix} = \vec{i} \begin{vmatrix} A_y & A_z \\ B_y & B_z \end{vmatrix} - \vec{j} \begin{vmatrix} A_x & A_z \\ B_x & B_z \end{vmatrix} + \vec{k} \begin{vmatrix} A_x & A_y \\ B_x & B_y \end{vmatrix}$$

determinant

-Vector operator del:  $\vec{\nabla} = \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z}$

For example, the linear momentum operator:

$$\vec{p} = -i\hbar \vec{\nabla}$$

-Derivative of the vector:  $\frac{d\vec{A}}{dt} = \vec{i} \frac{dA_x}{dt} + \vec{j} \frac{dA_y}{dt} + \vec{k} \frac{dA_z}{dt}$

-We can use vector notation to represent particle's coordinates:

$$\Psi(x_1, y_1, z_1, x_2, y_2, z_2) = \Psi\left(\vec{r}_1, \vec{r}_2\right)$$

Example:  $\vec{A}(4, -1, 3)$

$\vec{B}(-1, 2, 2)$

$|\vec{A}| = ?$

$\vec{C} = \vec{A} + \vec{B} = ?$

$D = \vec{A} \cdot \vec{B} = ?$

$\vec{E} = \vec{A} \times \vec{B} = ?$

$|\vec{A}| = \sqrt{16 + 1 + 25} = \sqrt{36} = 6$

$\vec{C}(3, 1, 7)$

$D = (-1) \cdot 4 + (-1) \cdot 2 + 3 \cdot 2 = -4 - 2 + 6 = 0$

$\vec{E} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 4 & -1 & 3 \\ -1 & 2 & 2 \end{vmatrix} = \vec{i} \begin{vmatrix} -1 & 3 \\ 2 & 2 \end{vmatrix} - \vec{j} \begin{vmatrix} 4 & 3 \\ -1 & 2 \end{vmatrix} + \vec{k} \begin{vmatrix} 4 & -1 \\ -1 & 2 \end{vmatrix}$

$\vec{E}(-8, -12, 7)$

### Back to Angular Momentum

Classical mechanics:  $\vec{L} = \vec{r} \times \vec{p}$

$\vec{L} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ x & y & z \\ p_x & p_y & p_z \end{vmatrix} = \vec{i} \begin{vmatrix} y & z \\ p_y & p_z \end{vmatrix} - \vec{j} \begin{vmatrix} x & z \\ p_x & p_z \end{vmatrix} + \vec{k} \begin{vmatrix} x & y \\ p_x & p_y \end{vmatrix}$

$L_x = y \cdot p_z - z \cdot p_y$

$L_y = -x \cdot p_z + z \cdot p_x$

$L_z = x \cdot p_y - y \cdot p_x$

$\vec{r} = x \cdot \vec{i} + y \cdot \vec{j} + z \cdot \vec{k}$

$\vec{p} = p_x \cdot \vec{i} + p_y \cdot \vec{j} + p_z \cdot \vec{k}$

## Quantum mechanics:

- Operators will be taken from classical physics:

Classical

$$L_x = y \cdot p_z - z \cdot p_y$$

$$L_y = -x \cdot p_z + z \cdot p_x$$

$$L_z = x \cdot p_y - y \cdot p_x$$



Quantum

$$\hat{L}_x = y \cdot \left( -i\hbar \frac{\partial}{\partial z} \right) - z \cdot \left( -i\hbar \frac{\partial}{\partial y} \right) = -i\hbar \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right)$$

$$\hat{L}_y = -i\hbar \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$\hat{L}_z = -i\hbar \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

-Can we obtain all three of them simultaneously?

We need to find out if they commute:

$$\left[ \hat{L}_x, \hat{L}_y \right] = ?; \left[ \hat{L}_y, \hat{L}_z \right] = ?; \left[ \hat{L}_z, \hat{L}_x \right] = ?$$

$$\left[ \hat{L}_x, \hat{L}_y \right] = \hat{L}_x \cdot \hat{L}_y - \hat{L}_y \cdot \hat{L}_x$$

$$\hat{L}_x \cdot \hat{L}_y f = -i\hbar \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \cdot \left( -i\hbar \left( z \frac{\partial f}{\partial x} - x \frac{\partial f}{\partial z} \right) \right) =$$

$$= -\hbar^2 \left( y \frac{\partial}{\partial z} \left( z \frac{\partial f}{\partial x} \right) - z \frac{\partial}{\partial y} \left( z \frac{\partial f}{\partial x} \right) - y \frac{\partial}{\partial z} \left( x \frac{\partial f}{\partial z} \right) + z \frac{\partial}{\partial y} \left( x \frac{\partial f}{\partial z} \right) \right) =$$

$$= -\hbar^2 \left( y \frac{\partial f}{\partial x} + y \cdot z \frac{\partial}{\partial z} \frac{\partial f}{\partial x} - z^2 \frac{\partial}{\partial y} \frac{\partial f}{\partial x} - y \cdot x \frac{\partial^2 f}{\partial z^2} + z \cdot x \frac{\partial}{\partial y} \frac{\partial f}{\partial z} \right)$$

$$\hat{L}_y \cdot \hat{L}_x f = -\hbar^2 \left( z \cdot y \frac{\partial^2 f}{\partial x \partial y} - x \cdot y \frac{\partial^2 f}{\partial z^2} + x \frac{\partial f}{\partial y} + x \cdot z \frac{\partial^2 f}{\partial z \partial y} \right)$$

$$\left[ \hat{L}_x, \hat{L}_y \right] = -\hbar^2 \left( y \frac{\partial f}{\partial x} - x \frac{\partial f}{\partial y} \right) = i\hbar \hat{L}_z$$

Similarly:

$$\left[ \hat{L}_y, \hat{L}_z \right] = i\hbar \hat{L}_x$$

$$\left[ \hat{L}_z, \hat{L}_x \right] = i\hbar \hat{L}_y$$

Thus, we cannot obtain all three components simultaneously.

What about magnitude?

Let's define magnitude operator as:

$$\hat{L}^2 = \left| \begin{matrix} \rightarrow \\ \hat{L} \end{matrix} \right|^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2$$

Some useful transformations with operators:

1.  $\left[ \hat{A}, \hat{B} \right] = -\left[ \hat{B}, \hat{A} \right]$
2.  $\left[ \hat{A}, \hat{B} + \hat{C} \right] = \left[ \hat{A}, \hat{B} \right] + \left[ \hat{A}, \hat{C} \right]$
3.  $\left[ \hat{A}, \hat{A}^n \right] = 0$
4.  $\left[ \hat{A}, \hat{B} \cdot \hat{C} \right] = \left[ \hat{A}, \hat{B} \right] \cdot \hat{C} + \hat{B} \cdot \left[ \hat{A}, \hat{C} \right]$

Now we need to find if magnitude operator commutes with operators for x, y and z component:

$$\left[ \hat{L}^2, \hat{L}_x \right] = ?$$

$$\begin{aligned}
 \overset{\text{Rule 1}}{\left[ \hat{L}^2, \hat{L}_x \right]} &= \left[ \hat{L}_x + \hat{L}_y + \hat{L}_z, \hat{L}_x \right] = \overset{\text{Rule 3}}{\left[ \hat{L}_x, \hat{L}_x \right]} + \left[ \hat{L}_y, \hat{L}_x \right] + \left[ \hat{L}_z, \hat{L}_x \right] = \left[ \hat{L}_y \cdot \hat{L}_y, \hat{L}_x \right] + \left[ \hat{L}_z \cdot \hat{L}_z, \hat{L}_x \right]
 \end{aligned}$$

$$\begin{aligned}
 \overset{\text{Rule 4}}{=} & \hat{L}_y \cdot \left[ \hat{L}_y, \hat{L}_x \right] + \left[ \hat{L}_y, \hat{L}_x \right] \cdot \hat{L}_y + \hat{L}_z \cdot \left[ \hat{L}_z, \hat{L}_x \right] + \left[ \hat{L}_z, \hat{L}_x \right] \cdot \hat{L}_z
 \end{aligned}$$

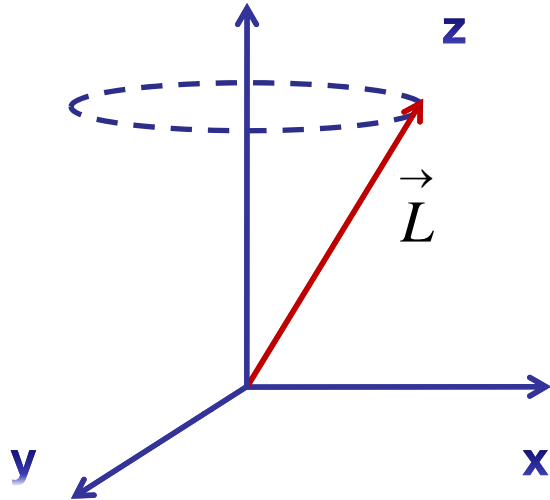
$$\begin{aligned}
 \overset{\text{Rule 2}}{=} & -\hat{L}_y \cdot \left[ \hat{L}_x, \hat{L}_y \right] - \left[ \hat{L}_x, \hat{L}_y \right] \cdot \hat{L}_y + \hat{L}_z \cdot \left[ \hat{L}_z, \hat{L}_x \right] + \left[ \hat{L}_z, \hat{L}_x \right] \cdot \hat{L}_z =
 \end{aligned}$$

$$= i\hbar(-\hat{L}_y \cdot \hat{L}_z - \hat{L}_z \cdot \hat{L}_y + \hat{L}_y \cdot \hat{L}_z + \hat{L}_z \cdot \hat{L}_y) = 0$$

Similarly:

$$\left[ \hat{L}^2, \hat{L}_x \right] = \left[ \hat{L}^2, \hat{L}_y \right] = \left[ \hat{L}^2, \hat{L}_z \right] = 0$$

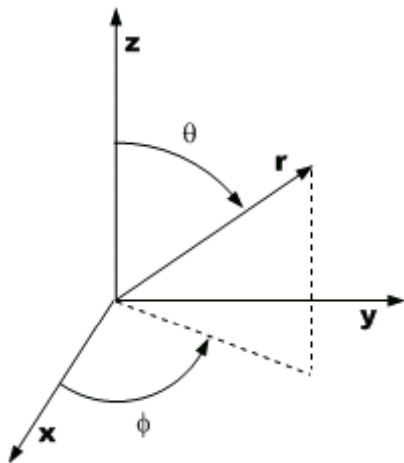
So, we can simultaneously determine the magnitude of orbital angular momentum of the particle and one of the components, say z:



*Vector  $L$  can lie anywhere on the surface of the cone.*

Even though we are using angular orbital momentum derived from classical physics, this does NOT mean that the electron is orbiting around nucleus.

Thus, we can use eigenvalue problem to obtain  $L_z$  and  $L^2$ . But, before doing that, we need to move to spherical coordinates in order to make partial derivatives separable).



$$x = r \sin \theta \cos \Phi$$

$$y = r \sin \theta \sin \Phi$$

$$z = r \cos \theta$$

$$r^2 = x^2 + y^2 + z^2$$

To convert the operators into spherical coordinates:

$$\overset{\text{cartesian}}{\downarrow} \quad \overset{\text{spherical}}{\downarrow}$$

$$\hat{L}_z = -i\hbar\left(x\frac{\partial}{\partial y} - y\frac{\partial}{\partial x}\right) = -i\hbar\frac{\partial}{\partial\Phi}$$

$$\hat{L}^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2 = -\hbar^2\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\Phi^2}\right)$$

Operators depend on three cartesian coordinates (x,y,z) but only two spherical coordinates ( $\theta, \Phi$ ).

Now, let's solve the eigenvalue problem:  $\hat{L}_z Y(\theta, \Phi) = b \cdot Y(\theta, \Phi)$

$$\hat{L}^2 Y(\theta, \Phi) = c \cdot Y(\theta, \Phi)$$

We try a separation of variables:  $Y(\theta, \Phi) = S(\theta) \cdot T(\Phi)$

Applying  $L_z$  operator:  $-i\hbar\frac{\partial}{\partial\Phi}[S(\theta) \cdot T(\Phi)] = b \cdot S(\theta) \cdot T(\Phi)$

$$-i\hbar S(\theta)\frac{\partial T(\Phi)}{\partial\Phi} = b \cdot S(\theta) \cdot T(\Phi)$$

$$\frac{dT(\Phi)}{T(\Phi)} = -\frac{b}{i\hbar} d\Phi$$

$$\ln(T(\Phi)) = \frac{ib\Phi}{\hbar} \Rightarrow T(\Phi) = A \cdot e^{\frac{ib\Phi}{\hbar}}$$

*Not a suitable wavefunction, it is not single-valued*

For T to be single-valued, we have the restriction:

$$T(\Phi + 2\pi) = T(\Phi)$$

$$A \cdot e^{\frac{ib\Phi}{\hbar}} \cdot e^{\frac{ib2\pi}{\hbar}} = A \cdot e^{\frac{ib\Phi}{\hbar}}$$

$$e^{\frac{ib2\pi}{\hbar}} = 1 \Rightarrow e^{i\alpha} = \cos \alpha + i \sin \alpha = 1 \longrightarrow \text{works for } \alpha = 2\pi m, \text{ where } m = 0, \pm 1, \pm 2, \dots$$

$$\frac{b2\pi}{\hbar} = 2\pi m \Rightarrow b = m\hbar \quad T(\Phi) = A \cdot e^{im\Phi} \quad \text{eigenvalues for } L_z \text{ are quantized}$$

$$m = 0, \pm 1, \pm 2, \dots$$

How do we get A?  
By normalization.

$$\int_0^{2\pi} |T(\Phi)|^2 d\Phi = 1$$

$$\int_0^{2\pi} (A \cdot e^{im\Phi}) \cdot (A \cdot e^{im\Phi})^* d\Phi = A^2 \int_0^{2\pi} d\Phi = A^2 2\pi$$

$$A = \frac{1}{\sqrt{2\pi}}$$

$$T(\Phi) = \sqrt{\frac{1}{2\pi}} \cdot e^{im\Phi}$$

Derivation of solution for eigenvalues of  $L^2$  operator is more complex, so we will present only the solution:

$$-\hbar^2 \left( \frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \Phi^2} \right) \left( S(\theta) \frac{1}{\sqrt{2\pi}} e^{im\Phi} \right) = c S(\theta) \frac{1}{\sqrt{2\pi}} e^{im\Phi}$$



$$S(\theta) = \sin^{|m|} \theta \sum_{j=1,3,\dots}^{l-|m|} a_j \cos^j \theta$$

$$l=0,1,2,\dots$$

$$|m| \leq l$$

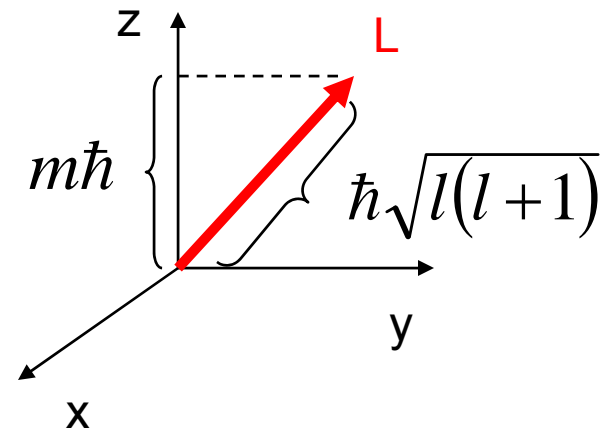
Possible values for m are:  $-l, -l+1, \dots, 0, 1, \dots, l-1, l$

$$c = l(l+1)\hbar^2$$

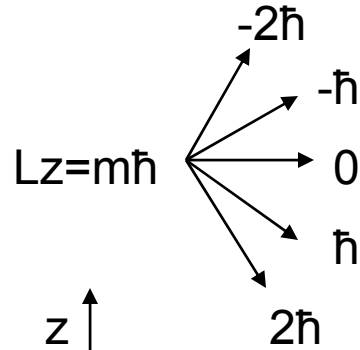
Thus, the final solutions are:

$$\hat{L}^2 Y = l(l+1)\hbar^2 Y \quad l = 0, 1, 2, \dots$$

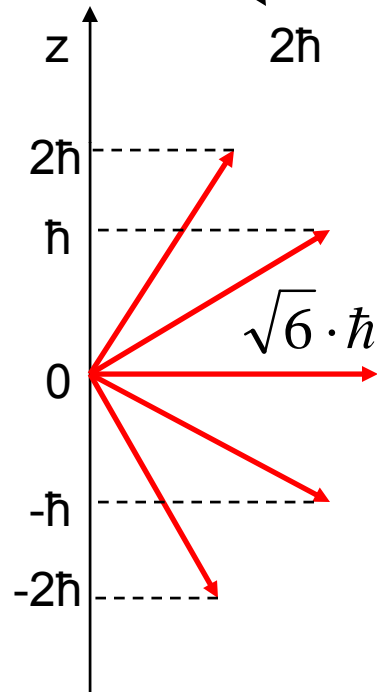
$$\hat{L}_z Y = m\hbar Y \quad m = -l, -l+1, \dots, 0, \dots, l-1, l$$



Example: if  $l=2$ , how many possible outcomes can we have?



$$|L| = \sqrt{6} \cdot \hbar$$



If  $\psi$  is an eigenfunction of  $L_z$  and  $L^2$  operators, how many outcomes?

*Only one outcome*

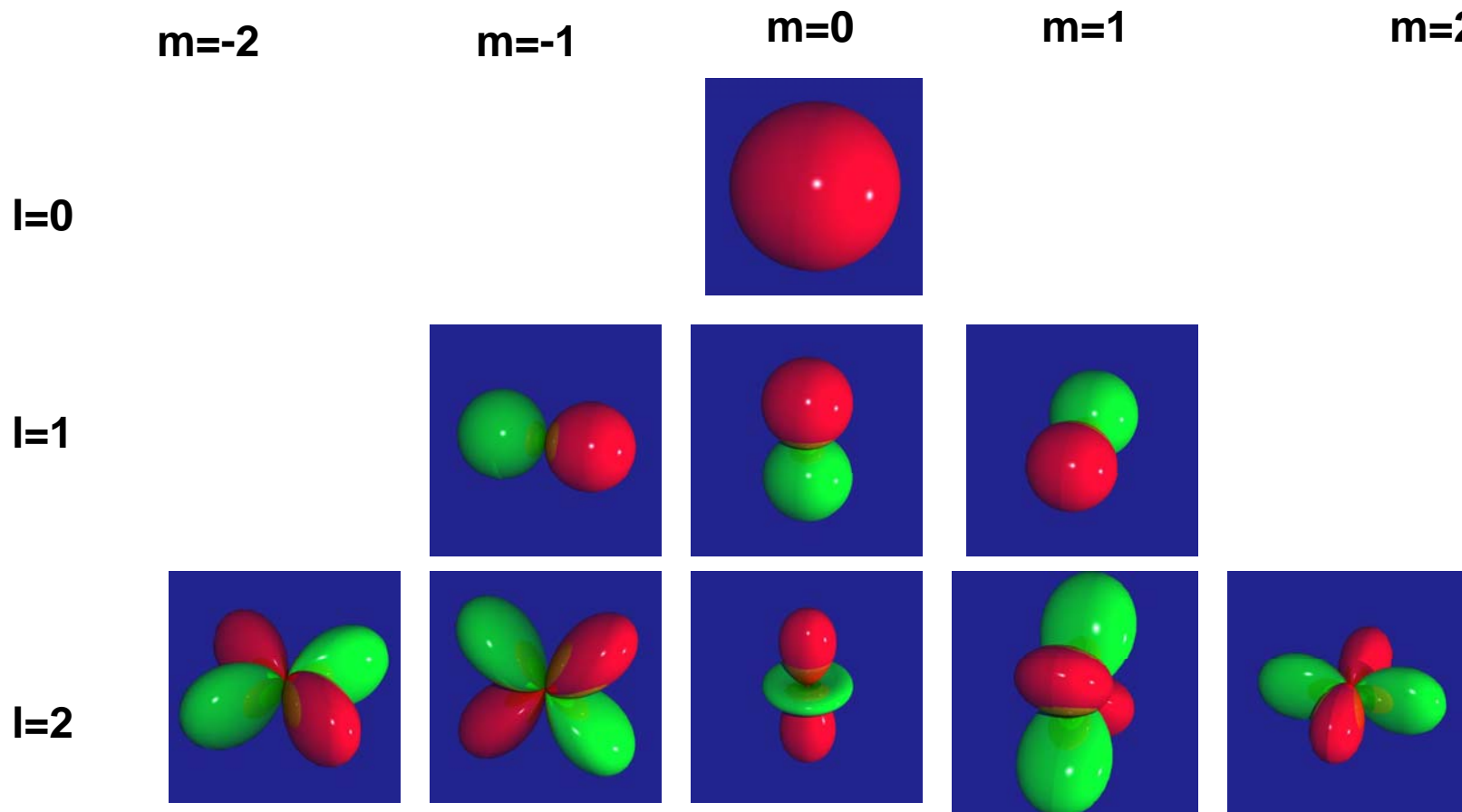
If  $\psi$  is not an eigenfunction of  $L_z$  and  $L^2$  operators, how many outcomes?

*Five outcomes*

Eigenfunctions  $Y$  are called spherical harmonics. They are complex functions. But, to visualize them, we can use the real combinations of spherical harmonics defined as:

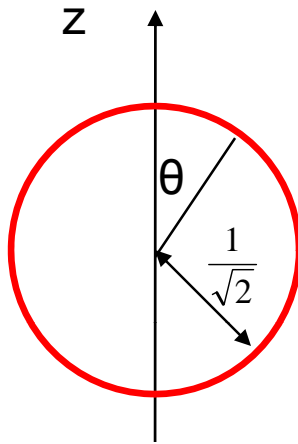
$$S_l^m = \frac{1}{\sqrt{2}} (Y_l^m + Y_l^{-m}) \qquad S_l^{-m} = \frac{1}{i\sqrt{2}} (Y_l^m - Y_l^{-m})$$

Here is how the combinations look like (the plot is made using *polar coordinates*, where the distance from the origin to a point on the graph is  $ST$ ):

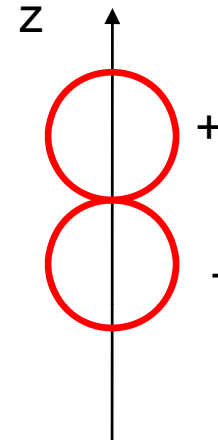


In polar coordinates, variables  $\theta$  and  $\Phi$  are expressed as angles with z-axis. For example:

$$S_{00} = \frac{1}{\sqrt{2}}$$



$$S_{10} = \frac{1}{2}\sqrt{6}\cos\theta$$



-The same spherical harmonics are obtained if we solved Schrodinger equation for a particle rotating on a sphere.

-To describe H-atom, we also treat it as a particle rotating on a sphere. The only difference is that the solution we obtained in this chapter involves a particle on a fixed distance  $r$  from the center, while in the case of H-atom, the distance between electron and a nucleus is not fixed. It is determined by Coulombic interactions between electron and a nucleus (attractive force that decays as  $1/r$ ). Thus, the final wavefunction  $\psi$  for H-atom will be a product of spherical harmonic  $Y$  and a radial function  $R(r)$ .